

Chapter 3 Rubber Elastic State

3.1 Introduction

Natural rubber : A latex from Hevea Brasiliensis

cf. rubber : material to remove marks from paper (지우개)

Elastomer : A polymer which exhibits rubber elastic properties

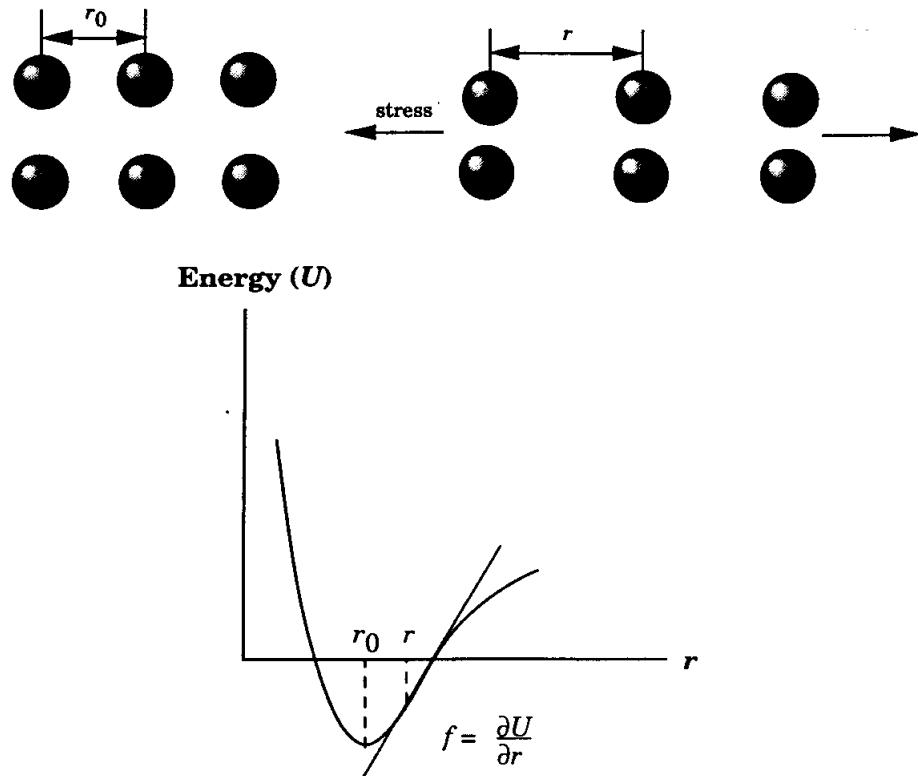
※ Rubber band experiment

100% strain by hanging a dead load

→ (heating) → suddenly decrease

- Metals or highly crystalline materials

→ Hookean elastic behavior (typically $\varepsilon < 0.2 \%$)



$$U \approx C(r - r_0)^2$$

C : constant

r_0 : equilibrium bond length

$$f = \frac{\partial U}{\partial r} = 2C(r - r_0)$$

f : elastic force

$$\sigma = \frac{2C}{A}(r - r_0) = E\varepsilon$$

$$E = 2Cr_0/A$$

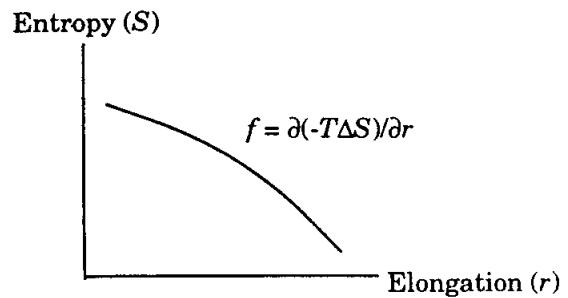
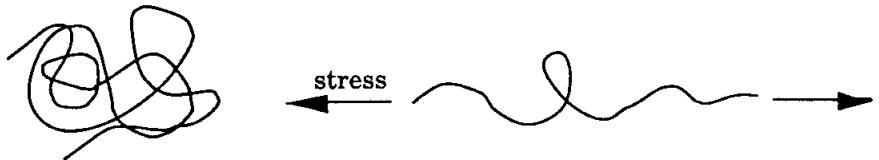
$$\varepsilon = \frac{r - r_0}{r_0}$$

\Rightarrow Hooke's law (energy-driven elasticity)

$\leftarrow U$ 가 최소되는 상태 선호

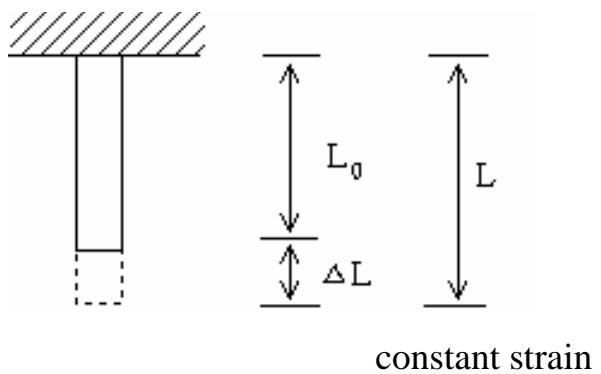
- **Rubbers** (entropy-driven elasticity) ← S 가 최대되는 상태 선호

The entropy of the rubber decreases on stretching.
(Due to changes in the conformational entropy)



$$f = \frac{\partial(-T\Delta S)}{\partial r}$$

3.2 Thermodynamics of rubber elasticity

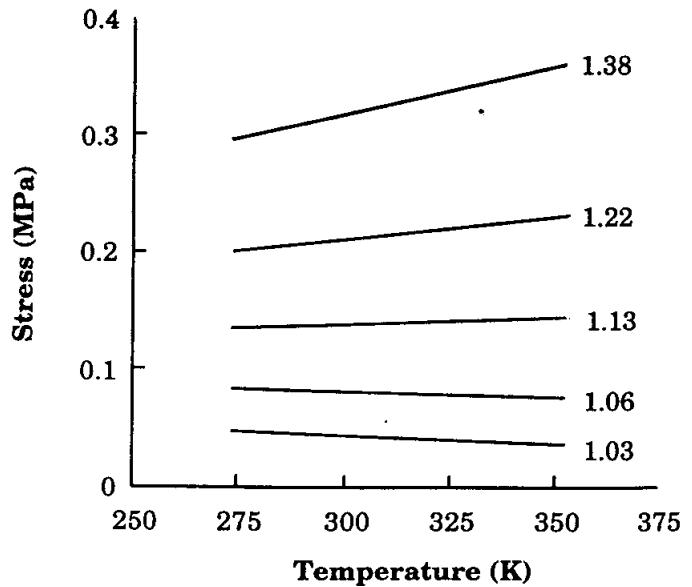


Measuring σ with various temperature

$$\lambda = L/L_0 < 1.1 \rightarrow \sigma \downarrow \text{ with } T \uparrow$$

$$\text{"} > 1.1 \rightarrow \sigma \uparrow \text{ with } T \uparrow$$

“thermo-elastic inversion”



From the thermodynamic laws,

$$dE = TdS - pdV + fdL \quad : \text{reversible process}$$

Internal energy change
Work done by deformation

$$\begin{aligned} G &= H - TS \\ &= E + pV - TS \quad (\text{미분하면}) \end{aligned}$$

$$dG = fdL + Vdp - SdT$$

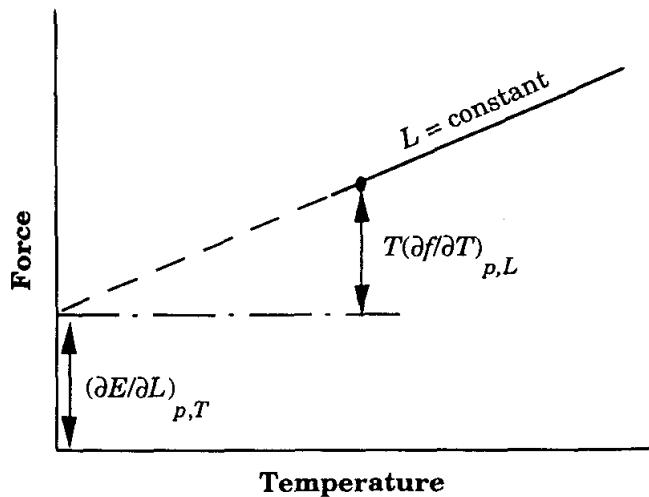
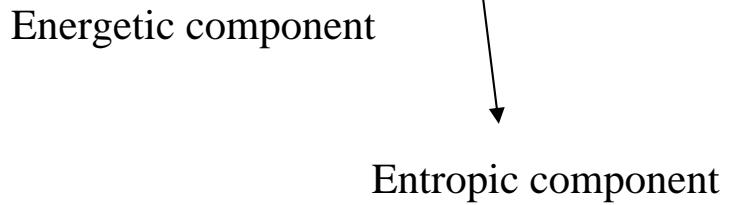
$$\left(\frac{\partial G}{\partial L} \right)_{T,p} = f, \quad \left(\frac{\partial G}{\partial T} \right)_{p,L} = -S$$

G : state function (exact differential) 즉, 미분순서에 상관없음

$$\left(\frac{\partial f}{\partial T}\right)_{L,p} = -\left(\frac{\partial S}{\partial L}\right)_{p,T} \quad : Maxwell\ relation$$

정리하면,

$$\Rightarrow f = \left(\frac{\partial E}{\partial L}\right)_{p,T} + T\left(\frac{\partial f}{\partial T}\right)_{p,L}$$



Analogous expression for constant V conditions

$$f = \left(\frac{\partial E}{\partial L} \right)_{V,T} + T \left(\frac{\partial f}{\partial T} \right)_{V,L}$$



Energetic force component

(f_e)

Entropic force component

(Natural rubber인 경우 80-85% 정도 차지)

Fraction of energetic force component

$$\frac{f_e}{f} = 1 - \frac{T}{f} \left(\frac{\partial f}{\partial T} \right)_{V,L}$$

f_e/f of PE : -0.42 \Rightarrow extended conformation requires lower energy

f_e/f of natural rubber : 0.15 - 0.2

Table 3.1 Energetic stress ratio of a few polymers

Polymer	Diluent	v_2^*	f_e/f
Polyethylene	none	1.00	-0.42
Polyethylene	$n\text{-C}_{30}\text{H}_{62}$	0.50	-0.64
Polyethylene	$n\text{-C}_{32}\text{H}_{66}$	0.30	-0.50
Natural rubber	none	1.00	0.17
Natural rubber	$n\text{-C}_{16}\text{H}_{34}$	0.34–0.98	0.18
Natural rubber	decalin	0.20	0.14
Poly(dimethyl siloxane)	none	1.00	0.25
Trans (1,4-polyisoprene)	none	1.00	-0.10
Trans (1,4-polyisoprene)	decalin	0.18	-0.20

Source: Mark (1984).

* Volume fraction of polymer in network.

3.3 Statistical theory of rubber elasticity

Affine network model vs. Phantom network model

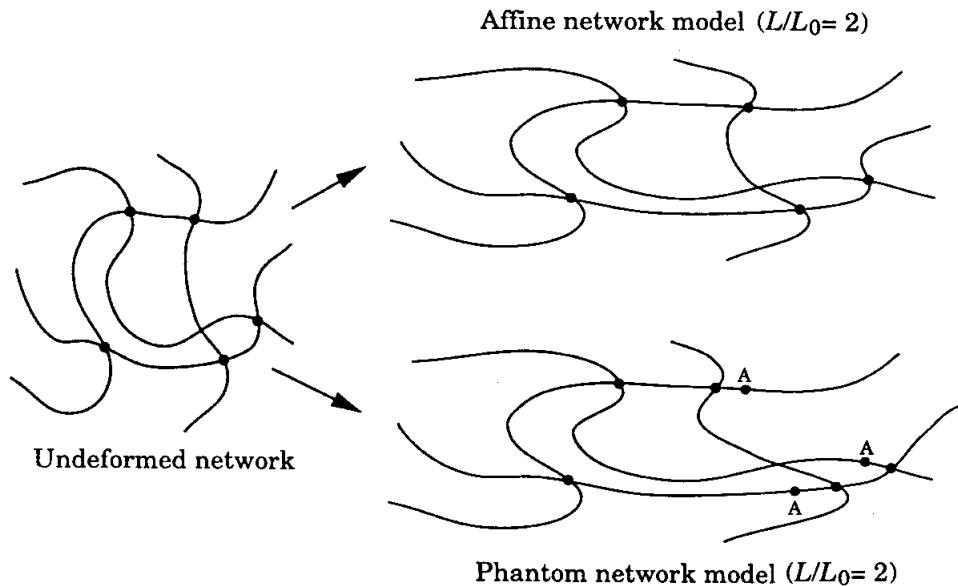


Figure 3.10 Schematic representation of the deformation of a network according to the affine network model (upper) and the phantom network model (lower). The points marked with an A indicate the position of the crosslinks assuming affine deformation.

Affine network model

Assumptions

- Chain segments are represented by Gaussian statistics
- S (entropy) of the network = Σ of S of the individual chains
- All different conformational states have the same energy
- Deformation is affine (균일 변형)
- Isotropic network at rest state
- Incompressible during deformation

Distribution of the end-to-end vectors

$$P(x, y, z) dx dy dz = \left(\frac{3}{2\pi \langle r^2 \rangle_0} \right)^{3/2} \exp\left(-3(x^2 + y^2 + z^2)/2\langle r^2 \rangle_0\right) dx dy dz$$

$\bar{r}_0 = (x_0, y_0, z_0)$: unstressed state

$\bar{r} = (x, y, z)$: stressed state

$$x = \lambda_1 x_0$$

$$y = \lambda_2 y_0 \quad \lambda_1, \lambda_2, \lambda_3 : \text{principal extension ratio (주 신장비)}$$

$$z = \lambda_3 z_0$$

$S = k \ln P$: Boltzmann's entropy relationship

$\Delta G = \Delta H - T\Delta S$ 의 관계식을 이용하면,



= 0 : All conformational states have the same energy

$$\rightarrow \Delta G = \frac{1}{2} nkT (\lambda_1^2 + \lambda_2^2 + \lambda_3^2 - 3)$$

$$nk = NR, R = N_A k$$

n : # of Gaussian chain segments

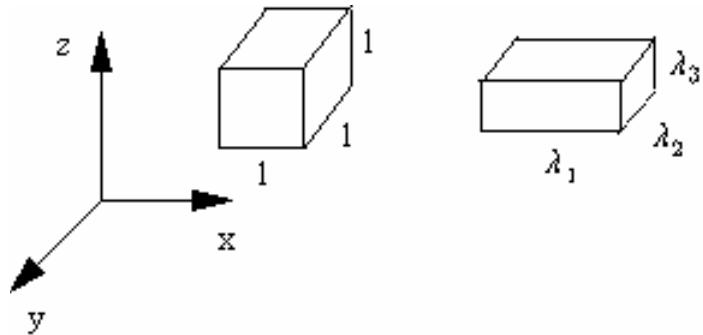
N : # of moles of Gaussian chain segments

* Stress-strain equation for a uniaxial stress

$$\lambda_1 = \lambda$$

$$\lambda_2 = \lambda_3 = \frac{1}{\sqrt{\lambda}} \leftarrow \lambda_1 \lambda_2 \lambda_3 = 1$$

$$\lambda = \frac{L}{L_0}$$



$$\begin{aligned} f &= \left(\frac{\partial(\Delta G)}{\partial L} \right)_{T,V} = \left(\frac{\partial(\Delta G)}{\partial \lambda} \right)_{T,V} \left(\frac{\partial \lambda}{\partial L} \right)_{T,V} \\ &= \frac{n k T}{L_0} \left(\lambda - \frac{1}{\lambda^2} \right) \end{aligned}$$

$$f = \sigma A = \sigma \left(\frac{A_0}{\lambda} \right) = \frac{NRT}{L_0} \left(\lambda - \frac{1}{\lambda^2} \right)$$

$$\therefore \sigma = \frac{NRT}{V_0} \left(\lambda^2 - \frac{1}{\lambda} \right)$$

$$\frac{N}{V_0} = \left(\frac{N\overline{M}_c}{V_0} \right) \left(\frac{1}{\overline{M}_c} \right) = \left(\frac{m_0}{V_0} \right) \left(\frac{1}{\overline{M}_c} \right) = \frac{\rho}{\overline{M}_c}$$

\overline{M}_c : No. avg. MW of the Gaussian chain segments

$$\therefore \sigma = \frac{\rho RT}{\overline{M}_c} \left(\lambda^2 - \frac{1}{\lambda} \right)$$

$$T \uparrow, \overline{M}_c \downarrow \rightarrow \sigma \uparrow$$

cf. *Phantom network model*

$$\sigma = \frac{\rho RT}{2\overline{M}_c} \left(\lambda^2 - \frac{1}{\lambda} \right) : \text{crosslink functionality } \geq 4 \text{인 경우}$$

3.5 Deviations from classical statistical theories

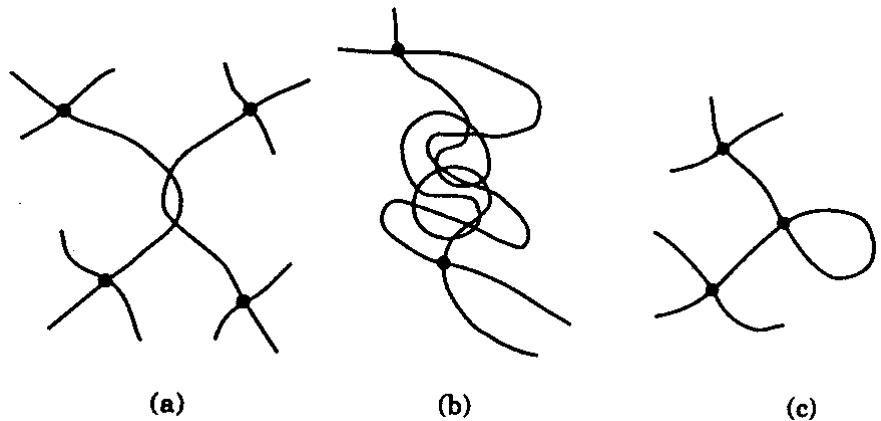
Loose chain ends do not contribute to the elastic force.

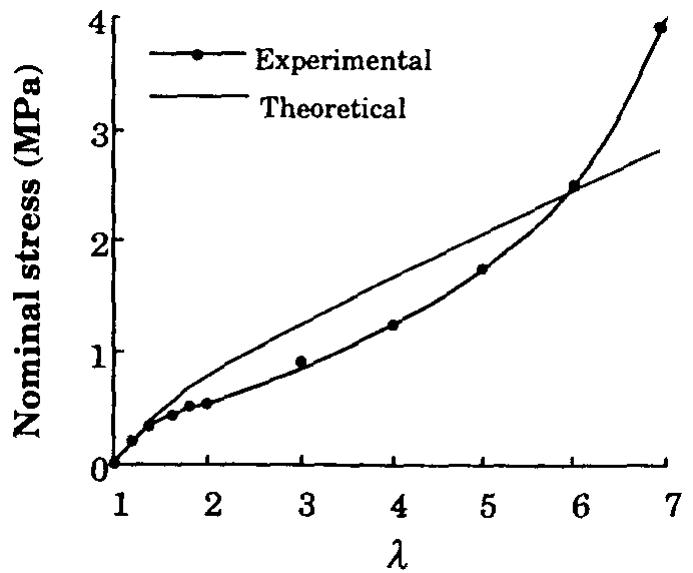
Modification of the stress-strain relationship

$$\sigma = \frac{\rho RT}{M_c} \left(1 - \frac{2\overline{M}_c}{M} \right) \left(\lambda^2 - \frac{1}{\lambda} \right)$$

Other types of network defect :

- (a) permanent physical crosslinks
- (b) temporary physical crosslinks
- (c) intramolecular crosslink





Above $\lambda = 4$, positive curvature

Possible causes :

- 1) non-Gaussian behavior
- 2) crystallization

3.7 Theory of Mooney & Rivlin

(based on *continuum mechanics*)

- For a uniaxially stressed specimen,

Basis: Rubber is incompressible

Isotropic

Hookean behavior in simple shear

Mooney-Rivlin equation:

$$\sigma = 2 \left(C_1 + \frac{C_2}{\lambda} \right) \left(\lambda^2 - \frac{1}{\lambda} \right)$$

C_1, C_2 : material constants

- . Related to the looseness with which the crosslinks are imbedded in the structure.
- . Solvent content $\uparrow \rightarrow C_2 \downarrow$
- . $v_2 \leq 0.2 \rightarrow C_2 \approx 0$

