Chapter 7. Transport Phenomena of Nanoparticles

7.1 Drag force

(1) Stokes' law

* Drag force, F_D :

- net force exerted by the fluid on the spherical particle in the direction of flow

 C_{D} : drag coefficient

cf. For pipe flow

$$
\tau_w = f \frac{\rho_f U^2}{2} \rightarrow F_w = f(\pi D L) \frac{\rho_f U^2}{2}
$$

where f : Fanning friction factor

*
$$
C_D
$$
 vs. Re_p
\n ρ , μ : density and viscosity of fluid
\n- For $Re_p < 1$ (creeping flow region)
\n $F_D = 3\pi d_p \mu U$ Stokes' law
\n
$$
F_D = \left(\frac{24}{Re_p}\right) \left(\frac{\pi}{4} d_p^2\right) \frac{\rho_f U^2}{2}
$$
\n
$$
\therefore C_D = \frac{24}{Re_p}
$$
\n
$$
cf. for pipe flow \t f = \frac{16}{Re}
$$

 $\frac{1}{2}$

- For 500
$$
$C_D = -0.44$
For $1 < Re_p < 500$ $C_D = \frac{24}{Re_p} (1 + 0.15 Re_p^{0.687})$
$$

(2) Non-continuum Effect

* Mean-free path of fluid

$$
\lambda = \frac{1}{\sqrt{2}n_m \pi d_m^2}
$$

where n_m : number concentration of molecules

 d_m : diameter of molecules

For air at 1 atm and 25° C, $\lambda = 65.1$ nm

For water at 25^oC , $\lambda = ?$

* Particle-fluid interaction

Continuum regime transition regime free-molecule regime

* Knudsen number -Continuum regime $K_n \sim 0$ (<0.1) -Transition regime $K_n \sim 1$ (0.1~10) -Free-molecular regime $K_n \sim \infty$ (>10) $\begin{array}{cc} n & d \end{array}$ $K_n = \frac{\lambda}{d_p}$ $K_n \sim 0$

- Particles in water is always in continuum regime…

* Corrected drag force

$$
F_d = \frac{3\pi d_p \mu U}{C_C}
$$

where C_{c} : Cunnigham correction factor [$C_C = 1 + K_n [20514 + 0.8 \exp(-0.55/K_n)]$

In air at 1 atm and 25°C

₫, J4792	С,
0.01	22.7
0.05	5.06
0.1	2.91
1.0	1.168
10	1.017

- Particles in water do not need noncontinuum correction…
Dynamic shape factors of powders

(3) Nonspherical particles χ * Shape correction factor, $d_{\,\nu} \mu U$ $F_{\scriptscriptstyle D}$ $\pi d_{\nu}\mu$ i χ \sim 3 $=\frac{1}{3\pi d_v}$

7.2 Migration in Gravitational Force Field

For the particle suspending in the fluid

- Force balance
$$
F_D = F_g - F_B
$$

\n
$$
\therefore C_D \left(\frac{\pi}{4} d^2 \rho \right) \frac{\rho_f U^2}{2} = \frac{\pi}{6} (\rho_p - \rho_f) d^3 \rho g
$$
\n
$$
\therefore U_T = \left[\frac{4}{3} \frac{gd \rho}{C_D} \left(\frac{\rho_p - \rho_f}{\rho_f} \right) \right]^{1/2}
$$
\nTerminal setting velocity

- For Stokes' law regime

^aFor unit density particle in air at 1atm and 25ºC

 Δ

* Dynamic equivalent diameters – calculated in air

Irregular particles and its equivalent spheres

In general

* Migration velocity

$$
F_{ext} = F_D \bigg(= \frac{3\pi\mu d_p}{C_c} U_m \bigg)
$$

where U_m : migration or drift velocity in the fields Note gravitational migration velocity: terminal settling velocity, U_T * Number flux by migration

$$
\vec{J}=n\vec{U}_m
$$

where *n*: particle number concentration

Centrifugal migrationAcceleration of centrifugation>>g 2 $1 - \frac{\rho_f}{\rho_p} \frac{U^2}{r} = m_p \left(1 - \frac{\rho_f}{\rho_p} \right) r \omega$ ρ ρ $\frac{\rho_f}{\rho} \left| \frac{U^2}{r} \right| = m_p \left(1 - \frac{\rho_f}{\rho} \right) r$ $\, U \,$ $F_c = m_p \left(1 - \frac{F}{\rho_p}\right) \frac{F}{r} = m_p \left(1 - \frac{F}{\rho_p}\right)$ $\frac{p}{p} \left(1 - \frac{\mu_f}{\rho}\right)$ t p $\sum_{c} = m_p \left(1 - \frac{P_f}{Q} \right)$ $= m_p \left(1 - \frac{\rho_f}{\rho_p}\right) \frac{U^2}{I} = m_p \left(1 - \frac{\rho_f}{\rho_p}\right)$ ($\left(\!\rho_{_{P}}-\rho_{_{f}}\right)$ r $dU_{cf} = \frac{(\rho_p - \rho_f) d^2 p U^2 C_c}{18 \mu r}$

Electrical Migration

$$
\vec{F} = q\vec{E} = n_e eE
$$

where q: charge of particles E : strength of electric field e: charge of electron (elementary unit of charge) n_e : number of the units p $e = \frac{e^{i\omega} - e^{i\omega}}{3\pi i d}$ $U_{\parallel} = \frac{n_e eEC}{\sqrt{2\pi}}$ $\frac{1}{2}$ 3 $\pi\mu$ $\therefore U_e =$

* Charging of particles

-Applied to electrostatic precipitation 7.3 Electrical Migration

$$
\vec{F} = q\vec{E} = n_e eE
$$

where *q*: charge of particles
E: strength of electric field
e: charge of electron (elementary unit of charge)
 n_e : number of the units

$$
\therefore U_e = \frac{n_e e E C_c}{3\pi \mu d_p}
$$
 Electrical migration velocity

-Electrical mobility

$$
Z_i = \frac{U_0}{E} = \frac{n_e e C_c}{3\pi\mu}
$$

-Applied to electrostatic precipitation in gas

- * Charging of particles in gas (Later for the case of liquid)
- Direct ionization
- -Static electrification: electrolyte, contact, spray, tribo, flame)
- -Collision with ions or ion cluster

* Diffusion charging

-Collision by ions and charged particles in Brrownian motion

$$
n(t) = \frac{d_p kT}{2K_E e^2} \ln\left[1 + \frac{\pi K_E d_p \overline{c_i} e^2 N_i t}{2kT}\right]
$$

Where $\,c_{_i}$: mean thermal speed of the ions(=240m/s at SC)

 N_i : ion concentration K_{E} : proportionality factor depending on unit used... K_{E} = $\frac{1}{4}$ \mathbf{u} 1 πε $K_{E} =$

* Field charging

-Charging by unipolar ions in the presence of a strong electric field

$$
n(t) = \left(\frac{3\varepsilon}{\varepsilon + 2}\right) \left(\frac{Ed_p^2}{4K_E e}\right) \left(\frac{\pi K_E e Z_i N_i t}{1 + \pi K_E e Z_i N_i t}\right)
$$

Where $\varepsilon_{\vec r}$ relative permissibility of the particle Z_i : mobility of ions(=0.00015m2/V s

-Saturated after sufficient time…

* Charge limit

-By electron ejection from mutual repulsion on the surface

$$
n_{\text{max}} = \frac{d_p^2 E_L}{4K_E e}
$$

where E_L : surface field strength required for spontaneous emission of electrons(=9.0 \times 10 8 V/m)

-Rayleigh limit

If mutual repulsion > surface tension force for liquid droplets

$$
n_{\max} = \left(\frac{2\pi\sigma d_p^{3}}{K_E e^2}\right)^{1/2}
$$

* bipolar charging: Kr_{85}

Electrophoresis

- -Movement of nanoparticles in liquid medium
- * Zeta potential : potential at the slip plane*
	- -Plane that separates the tightly bound liquid layer from the rest of liquid
	- -~Stern layer
	- determines the stability of colloidal dispersion or a sol
	- requires $>$ 25mV for the stability

$$
\zeta = \frac{q}{2\pi\varepsilon_0\varepsilon_r d_p (1 + \kappa d_p / 2)}
$$

* Electrical migration velocity

$$
U_E = \frac{2\varepsilon_0 \varepsilon_r \zeta E}{3\pi\mu}
$$

* Electrical mobility

$$
B_E = \frac{U_0}{E} = \frac{2\varepsilon_0 \varepsilon_r \zeta}{3\pi\mu}
$$

7.3 Migration by Interaction with Fluids

(1) Diffusion

* Brownian motion

: Random wiggling motion of particles by collision offluid molecules on them

***Brownian Diffusion:**

Particle migration due to concentration gradient by Brownian motion

 $\vec{J}=-D_p\vec{\nabla}n$

Fick's law

where D_{p} : diffusion coefficient of particles, cm²/s

n : particle concentration by number

cf. Diffusion of molecules

* Coefficient of Diffusion

$$
D_p = \frac{kTC_c}{3\pi\mu d_p}
$$

Diffusion Coefficient of Unit-density sphere at 20°C in air

cf. Liquid diffusivity 10⁻⁵cm²/s

- Mass balance for the cube in the fluid
\n
$$
\frac{\partial n}{\partial t} \Delta x \Delta y \Delta z = [J_x - (J_x + \Delta J_x)] \Delta y \Delta z
$$
\n
$$
+ [J_y - (J_y + \Delta J_y)] \Delta x \Delta z
$$
\n
$$
+ [J_z - (J_z + \Delta J_z)] \Delta x \Delta y
$$
\n
$$
+ [J_z - (J_z + \Delta J_z)] \Delta x \Delta y
$$
\n
$$
= -\Delta J_x \Delta y \Delta z - \Delta J_y \Delta x \Delta z - \Delta J_z \Delta x \Delta y
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* One-dimensional diffusion from the origin

$$
\frac{\partial n}{\partial t} = \vec{\nabla} \cdot \left(D_p \vec{\nabla} n \right) = D_p \nabla^2 n \longrightarrow \frac{\partial n}{\partial t} = D_p \frac{\partial^2 n}{\partial x^2}
$$

$$
At t=0 n=0 for all x except x=0
$$

At x=0, n=n₀ for all t and $\left(\frac{\partial n}{\partial x}\right)_{x=0} = 0$

The solution is :

$$
n(\eta) = 1 - \frac{2}{\pi} \int_{0}^{\eta} e^{-\eta^2} d\eta = 1 - erf(\eta) \equiv erf(\eta)
$$

Where
$$
\eta = \frac{z}{2\sqrt{D_p t}}
$$

Differentiating with respect to x

$$
n(x,t)\frac{dn(x,t)}{n_0} = \frac{1}{(4\pi D_p t)^{1/2}} \exp\left(\frac{-x^2}{4D_p t}\right) dx
$$

Normal distribution with respect to x-axis

-*Mean displacement:* $x = 0$

-Standard deviation or root-mean square displacement

$$
\sigma = x_{\rm rms} = \sqrt{2D_p t}
$$

represent particle movement (displacement) by diffusion

Net displacement in 1s due to Brownian motion and gravity for standard-density spheres at standard conditions

Particle diameter, µm	x_{rms} in 1s(m)	Settling in 1s(m)	X_{rms} / X_{sett}
0.01	3.3×10^{-4}	6.9×10^{-8}	<i>4800</i>
0. 1	3.7×10^{-5}	8.8×10^{-7}	42
1.0	7.4×10^{-6}	3.5×10^{-6}	0.21
	2.2×10^{-6}	3.1×10 ⁻³	7.1×10^{-4}

(2) Thermophoresis

-Examples of thermophoresis

Dust free surface on radiator or wall near it

Movement cigarette smoke to cold wall or window

Spoiling of the surface of the cold wall

Scale formation on the cold side in the heat exchanger

* In free molecular regime

Waldmann and Schmidt(1966) $\langle F_{th}=-p\lambda d\rangle_p-\frac{1}{T}$ independent of d_p $\, T \,$ $F_{\mu} = -p\lambda d^2\omega$ $\vec{F}_{th} = -p \lambda d^2 P \frac{\vec{\nabla}}{q}$ $\, T \,$ $\, T \,$ $\, T \,$ $\, T \,$ $\vec{U}_{th} = -\frac{3v\nabla T}{\left(\frac{\pi\alpha}{\sqrt{2}}\right)} = \sim 0.55v\frac{\nabla}{\sqrt{2}}$ $\left(1+\frac{\pi\alpha}{8}\right)$ ∇ $\therefore U_{th} = \therefore \vec{U}_{th} = -\frac{3\vec{v}T}{\sqrt{3\pi}} = \sim 0.55\vec{v} \frac{\vec{\nabla}}{3}$ ν $\frac{1}{\pi \alpha}$ = 0.55 84 ¹ 3

* Correction for continuum fluid-particle interaction

Brock(1962) where $\, T \,$ dH $\, T \,$ $F_{\mu} = \frac{\mu}{\mu}$ $\mathcal{L} \mathcal{P} G$ $\frac{d}{dt} = \frac{2\rho_{G}T}{\rho_{G}}$ $\frac{\pi\mu}{2}$ 9 $-9\pi r^2$ ∇= ------------- $\vec{F} = -9\pi\mu^2 d_{p}H\vec{\nabla}$ I l $\overline{\bigwedge}$ $\bigg($ +++ $+ 6Kn$ $1 + 2\frac{k_G}{m} + 8.8Kn$ k_{\perp} kKn k_{\perp} kKnHp $\frac{G}{\sqrt{2}}$ $\frac{p}{2}$ $\frac{G}{\sqrt{2}}$ $1 + 2\frac{v}{k} + 8.8$ 4.4 $1 + 6$ 1 \sim $\, T \,$ $\, C \,$ H $\, T \,$ $U_{\mu} = \frac{C_{\mu} + C_{c} - C_{c}}{C_{c}}$ $\epsilon_{th} = \frac{-\frac{1}{2}\rho_{c}c_{c}}{2\rho_{c}}$ $\frac{\mu}{2}$ −3∇ $\therefore U_{th} = \frac{\cdot}{\cdot}$ $\vec{=}$ $\vec{ }$

Terminal settling and thermophoretic velocities in a temperaturegradient of 1ºC/cm at 293K

G

l

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l

 $\bigg)$

e.g. submicron particles in the upper atmosphere

e.g. tails of comet, laser-lift of particles

7.4 Inertial Motion and Impact of Particles

(1) Inertial motion

-For Stokesian particles

Momentum (force) balance for a single sphere

Neglecting buoyancy force

$$
\therefore \frac{\pi}{6} d^3{}_p \rho_p \frac{dU}{dt} = -\frac{3\pi \mu d_p U}{C_c}
$$

$$
m_p \frac{dU}{dt} = -F_D
$$

 $-t/$

 $\tau = \frac{\rho}{\sqrt{2\pi}}$

18

=

 $_{p}d_{p}^{2}C_{c}$ d

 $\, C \,$

Net displacement in 1s due to Brownian motion and gravity for standard-density spheres at standard conditions

Integrating twice \int $\bigg)$ \setminus $=U_{\alpha}\tau$ −− τ | $1-e^{-\tau t}$ $x = U_0 \tau \left(1 - e^{-\tau}\right)$ $\,$ 1e 0ν \vert \vert 2

Integrating once, defining relation time as $\tau = \frac{p}{18\mu}$

 U $=$ $U_0e^{-/\tau}$

$$
As \frac{t}{\tau} \to \infty, \qquad x \sim U_0 \tau = \frac{\rho_p d^2 p U_0 C_c}{18 \mu} \equiv s
$$

* out of Stokes' range

(2) Simiulitude Law for Impaction : Stokesian Particles

* Impaction: deposition by inertia

- For $Re \leq 1$

Force balance around a particle

Defining dimensionless variables where U , L : characteristic velocity and length of the system $\left(\overrightarrow{U_p}-\overrightarrow{U_f}\right)$ $p \cup p \quad \circ f$ $p \frac{p}{dt}$ d $\,$ $\,$ dt \rightarrow \rightarrow \rightarrow \rightarrow d $\, U \,$ $m_{\scriptscriptstyle -}-$ = $-3\pi\mu$ L $\,$ $U_1 \equiv \frac{U}{L}$, $U_{f1} \equiv \frac{U}{L}$ $\,$ U_{α} = $f_1 = \frac{\sigma_f}{L}$ and $\theta = \frac{\sigma_f}{L}$ tU, $U_{f1} \equiv \frac{J}{I}$ and $\theta \equiv \frac{V}{I}$ ≡3 $\frac{3}{1}$ Luv $\frac{1}{1}$ LdUπ

$$
\frac{\pi}{6}d_p^3 \frac{LdU_1}{(L/U)d\theta} = -3\pi\mu d_p \left(L\overrightarrow{U_1} - L\overrightarrow{U_{f1}}\right)
$$
(subscript p: omitted)

$$
\frac{\rho_p d_p^2 U}{18\mu L} \frac{d\overrightarrow{U_1}}{d\theta} = -\left(\overrightarrow{U_1} - \overrightarrow{U_{f1}}\right)
$$

Define Stokes number

,

$$
St \equiv \frac{\rho_p d^2_{p} U}{18 \mu L} = \frac{\tau U}{L} \equiv \frac{\text{particle persistence}}{\text{size of obstacle}}
$$

:. $St \frac{dU_1}{d\theta} = -(\overrightarrow{U_1} - \overrightarrow{U_{f1}})$

Or in terms of displacement

$$
St \frac{d^2 \vec{r}_1}{d\theta^2} + \frac{d \vec{r}_1}{d\theta} = \overrightarrow{U_{f1}}
$$
\nwhere $\vec{r}_1 = \frac{\vec{r}}{L}$, \vec{r} : displacement vector
\n- The solution gives $\overrightarrow{r_1(\theta)}$, $\overrightarrow{r(t)}$ or where the particle is at time t...\n* Particle trajectory
\n $\therefore \overrightarrow{r_1(\theta)} = f(St, Re, R)$

-Two particle impaction regimes are similar

when the geometric, hydrodynamic and particle trajectories are the same...

-Applications

Cyclone, particle impactor, filter